

# Gauss' Law and Applications

Physics 2415 Lecture 5

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# Today's Topics

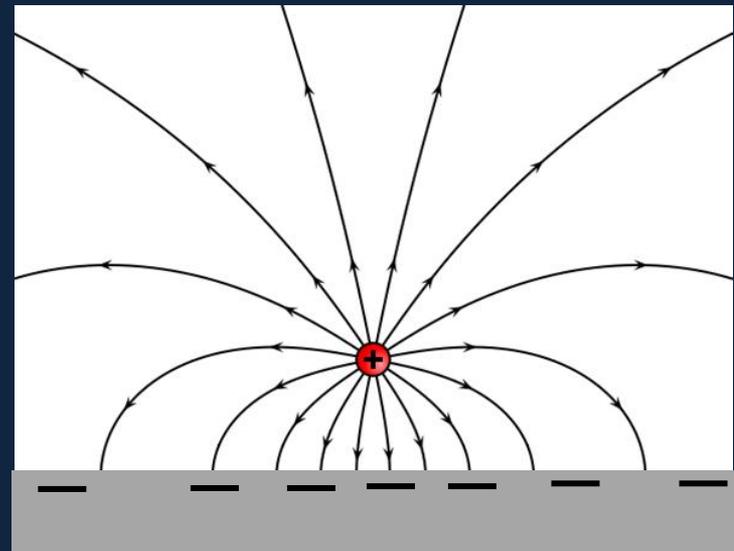
- Gauss' Law: where it came from—review
- Gauss' Law for Systems with Spherical Symmetry
- Gauss' Law for Cylindrical Systems: Coaxial Cable
- Gauss' Law for Flat Plates

# Clicker Question

- A charge  $+Q$  is placed a small distance  $d$  from a large flat **conducting** surface.
- **Describe the electric field lines:** close to the charge, they point radially outwards from the charge, but as they approach the conducting plane:
  - A. they bend away from it.
  - B. they reach it and just stop.
  - C. they curve around to meet the plane at right angles.

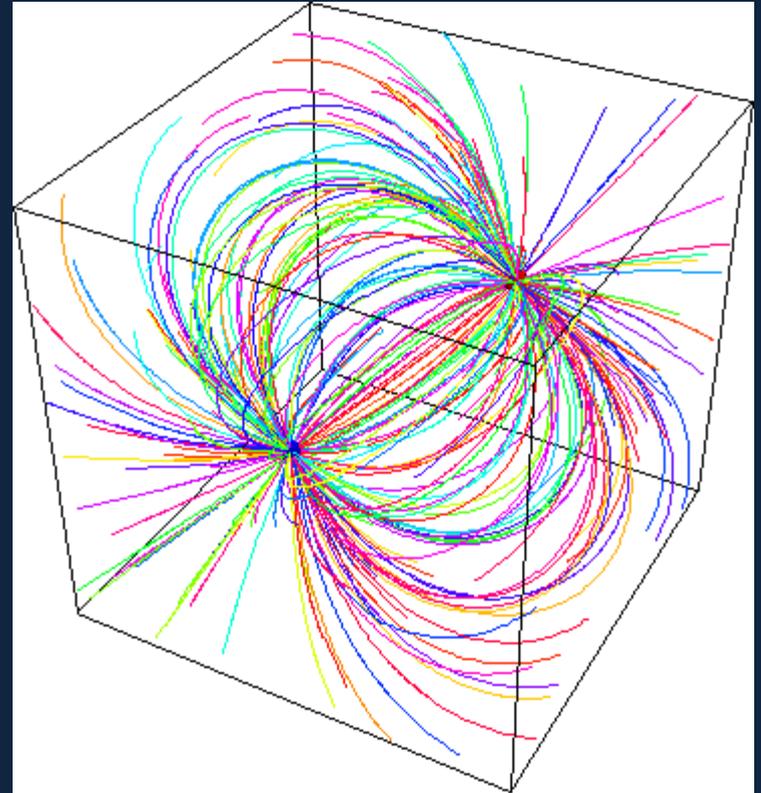
# Clicker Answer

- Field lines must always meet a conductor at right angles in electrostatics.
- Physically, the positive charge has attracted negative charges in the conductor to gather in the area under it. They repel each other, so are rather spread out.



# Dipole Field Lines in 3D

- There's an analogy with **flow of an incompressible fluid**: imagine fluid emerging from a source at the positive charge, draining into a sink at the negative charge.
- **The electric field lines are like stream lines**, showing fluid velocity direction at each point.
- Check out the applets at <http://www.falstad.com/vector2de/> !

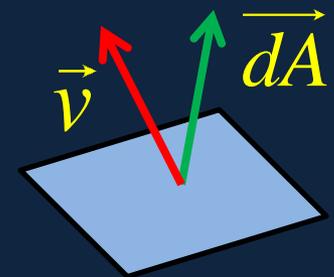
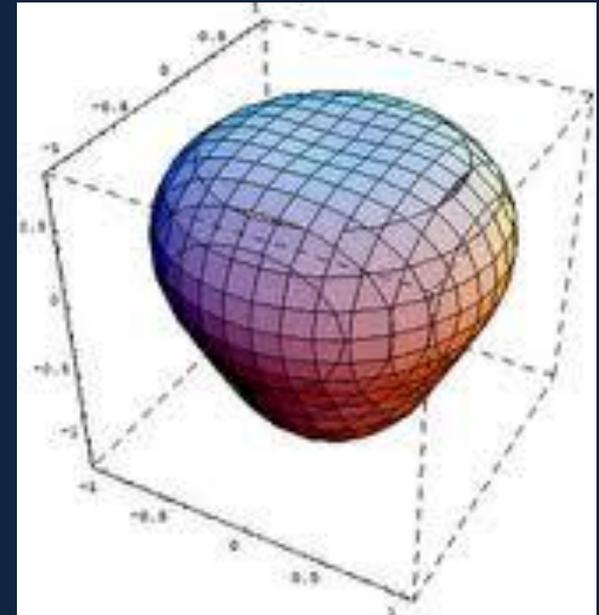


# Velocity Field for a Steady Source in 3D

- Imagine you're filling a deep pool, with a hose and its end, deep in the water, is a porous ball so the water flows out equally in all directions.
- Now picture the flow through a **spherical fishnet**, **centered on the source**, and far smaller than the pool size.
- Now think of a **second** spherical net, twice the radius of the first, so 4x the surface area. In steady flow, total water flow across the two spheres is the same: so  $v \propto 1/r^2$ .
- **This velocity field is identical to the electric field from a positive charge!**

# Total Flow through any Surface

- But how do we *quantify* the fluid flow through such a net?
- We do it **one fishnet hole at a time**: unlike the sphere, the **flow velocity is no longer always perpendicular to the area**.
- We represent each fishnet hole by a vector  $\vec{dA}$ , magnitude equal to its (small) area, direction perpendicular outwards. Flow through hole is  $\vec{v} \cdot \vec{dA}$
- The total outward flow is  $\int_{net} \vec{v} \cdot \vec{dA}$ .



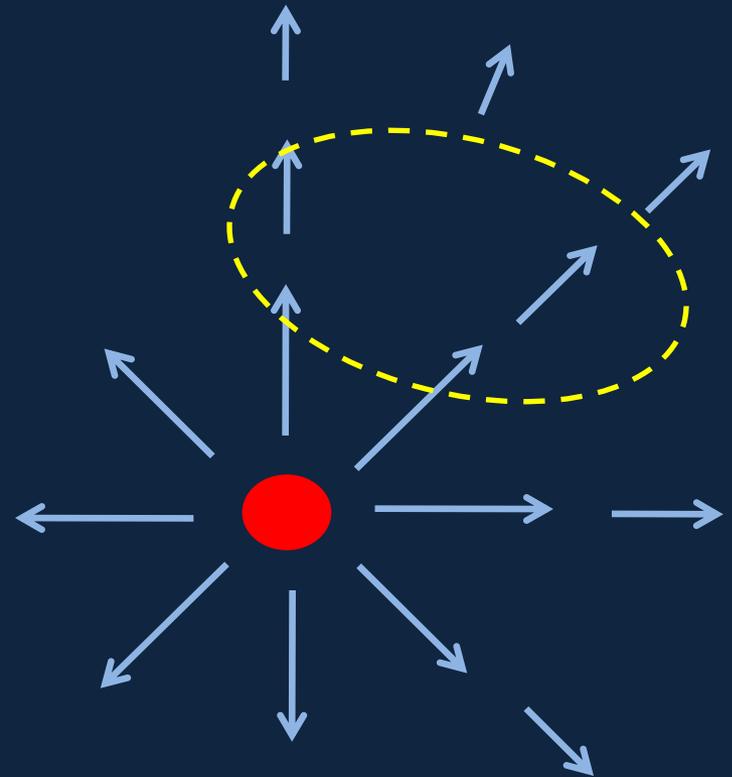
The component of  $\vec{v}$  perp. to the surface is  $v \cos \theta$ .

# Gauss's Law

- For incompressible fluid in steady outward flow from a source, the flow rate across any surface enclosing the source  $\int \vec{v} \cdot \vec{dA}$  is the same.
- The electric field from a point charge is identical to this fluid velocity field—it points outward and goes down as  $1/r^2$ .
- It follows that for the electric field  $\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{Q\hat{r}}{r^2}$  for any surface enclosing the charge  $\int \vec{E} \cdot \vec{dA} = \text{const.} = Q / \epsilon_0$  (the value for a sphere).

# What about a Closed Surface that *Doesn't* Include the Charge?

- The **yellow** dotted line represents some fixed closed surface.
- Think of the fluid picture: in steady flow, it goes in one side, out the other. The *net* flow across the surface must be zero—it can't pile up inside.
- By analogy,  $\int \vec{E} \cdot d\vec{A} = 0$  if the charge is outside.



# What about More than One Charge?

- Remember the **Principle of Superposition**: the electric field can always be written as a linear sum of contributions from individual point charges:

$$\vec{E} = \vec{E}_1 + \vec{E}_2 + \vec{E}_3 + \dots \text{ from } Q_1, Q_2, Q_3 \dots$$

and so

$$\int \vec{E} \cdot d\vec{A} = \int \vec{E}_1 \cdot d\vec{A} + \int \vec{E}_2 \cdot d\vec{A} + \int \vec{E}_3 \cdot d\vec{A} + \dots$$

will have a contribution  $Q_i / \epsilon_0$  from each charge inside the surface—this is **Gauss' Law**.

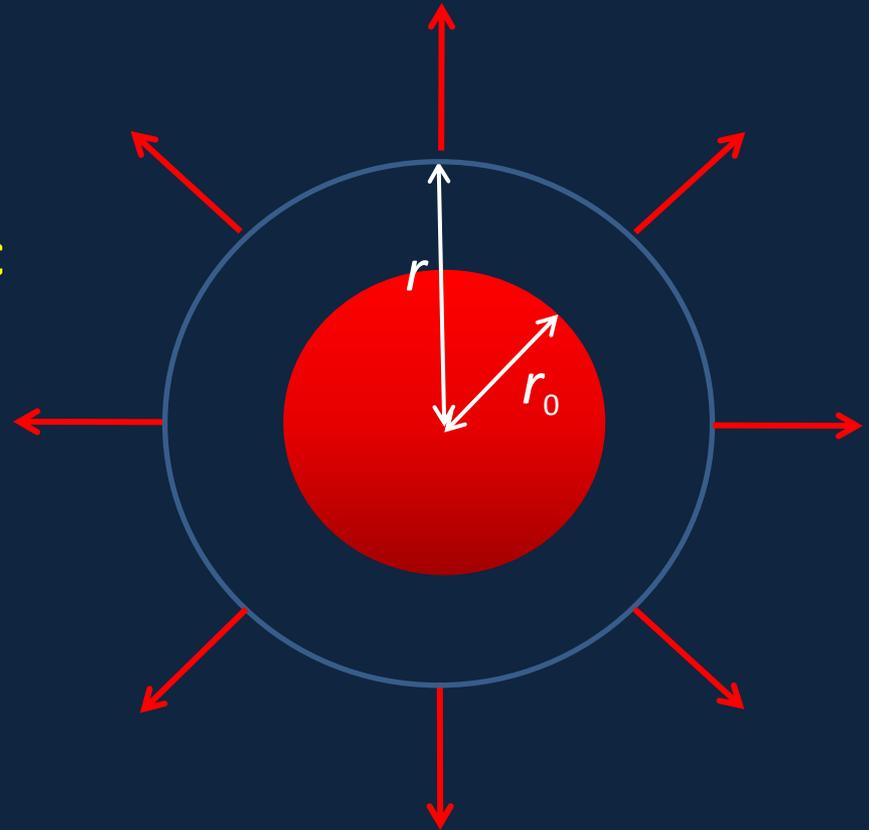
# Gauss' Law

- The integral of the total electric field flux out of a **closed surface** is equal to the **total charge  $Q$  inside the surface** divided by  $\epsilon_0$ :

$$\int_S \vec{E} \cdot d\vec{A} = \frac{Q}{\epsilon_0}$$

# Spherical Symmetry

- First, a **uniform spherical shell**, radius  $r_0$ , of positive charge.
- The perfect spherical symmetry means the **electric field outside**, at a distance  $r$  from the center, **must point radially outwards**. (rotating the sphere doesn't change anything, but *would* change a field pointing any other way.)

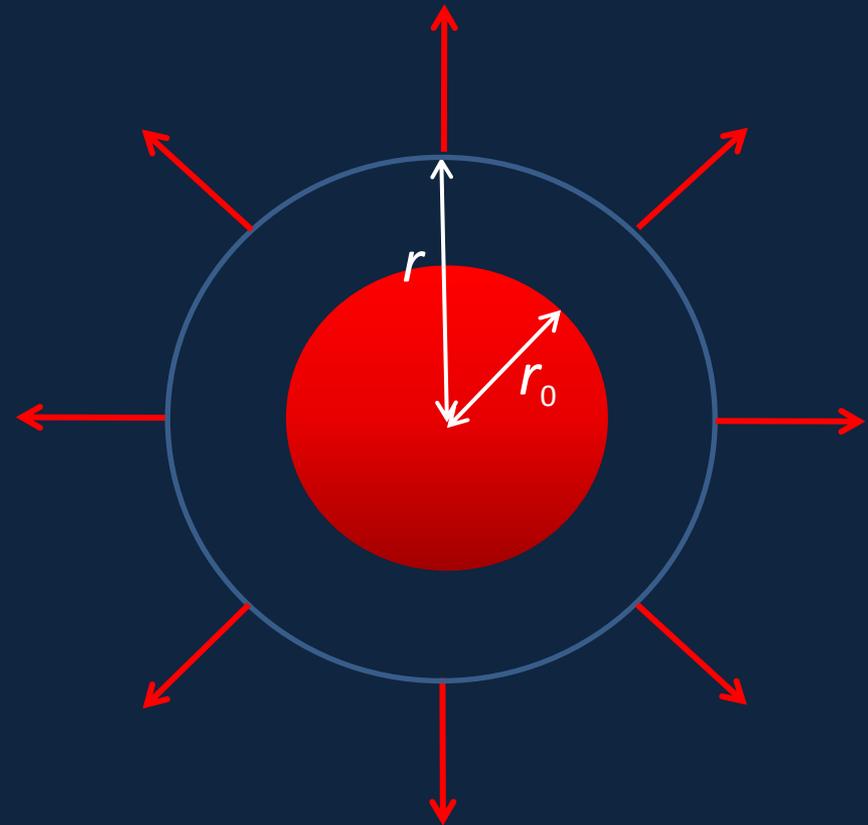


# Spherical Symmetry

- The blue circle represents a spherical surface of radius  $r$ , concentric with the shell of charge.
- For this enclosing surface, Gauss' Law  $\int_s \vec{E} \cdot d\vec{A} = \frac{Q}{\epsilon_0}$  becomes

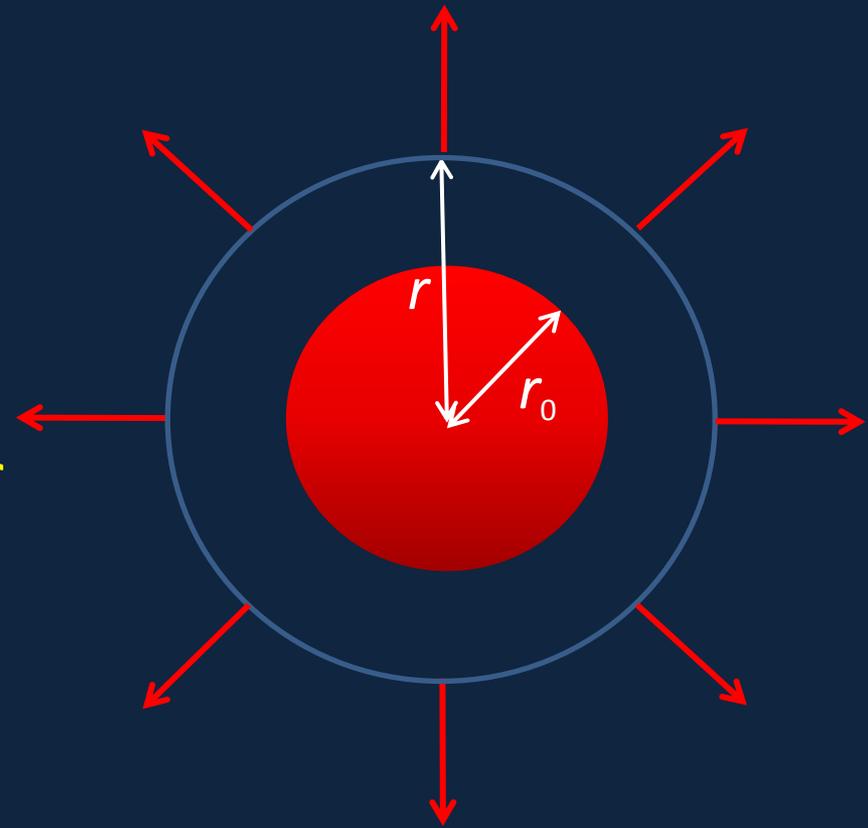
$$4\pi r^2 E(r) = Q / \epsilon_0,$$

$$E(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} = \frac{kQ}{r^2}$$



# Spherical Symmetry

- Gauss' Law easily shows that the electric field from a uniform shell of charge is the same outside the shell as if all the charge were concentrated at a point charge at the center of the sphere. This is difficult to derive using Coulomb's Law!

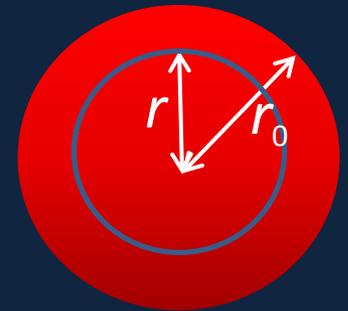


# Field *Inside* a Hollow Shell of Charge

- Now let's take the **enclosing surface** inside the hollow shell of charge.
- Gauss' Law is now

$$\int_S \vec{E} \cdot d\vec{A} = \frac{Q}{\epsilon_0} = 0$$

- Because there is no charge inside the shell, it's all on the surface.
- The spherical symmetry tells us the field inside the shell is exactly zero—again, not so simple from Coulomb's Law.



# Field Outside a Solid Sphere of Charge

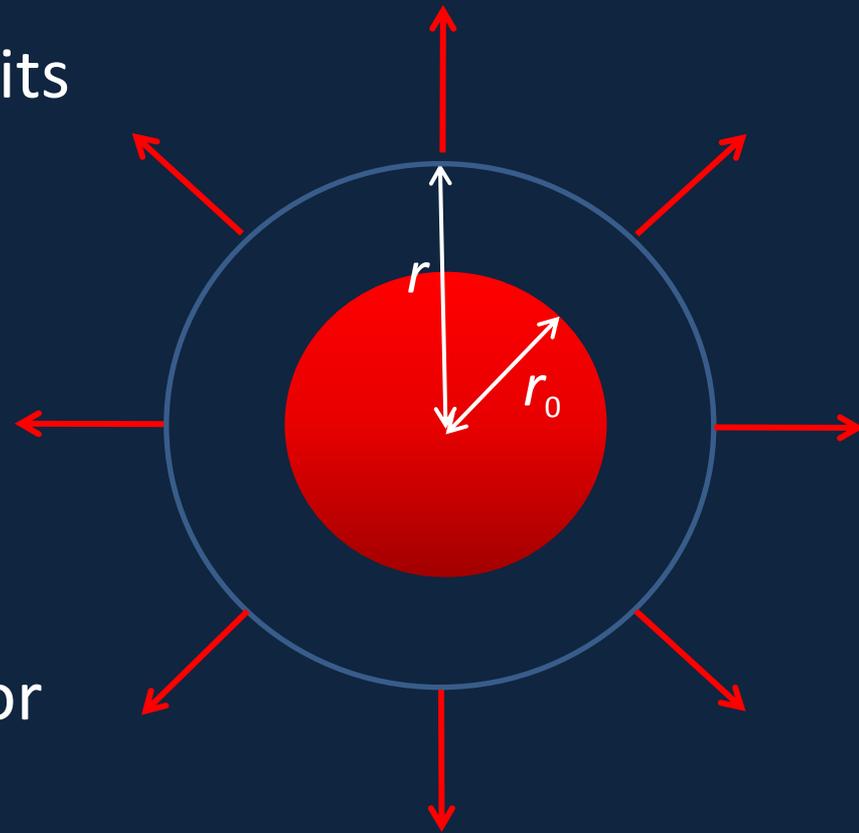
- Assume we have a sphere of insulator with total charge  $Q$  distributed uniformly through its volume.

- The field outside is again

$$E(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$

from the spherical symmetry.

- Note: Gauss' Law also works for *gravitation*—and **this is the result for a solid sphere of mass.**



# Field *Inside* a Solid Sphere of Charge

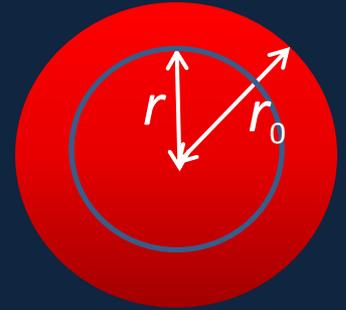
- Now let's take the enclosing surface inside the solid sphere of charge.
- Gauss' Law is now

$$\int_S \vec{E} \cdot d\vec{A} = \frac{Q_{\text{inside}}}{\epsilon_0} = Q \times \frac{\text{volume inside } S}{\text{total volume}} = \frac{Q}{\epsilon_0} \frac{r^3}{r_0^3}$$

- From this, since  $\int_S \vec{E} \cdot d\vec{A} = 4\pi r^2 E$ ,

$$E = \frac{Q}{4\pi\epsilon_0} \cdot \frac{r}{r_0^3}$$

so the electric field strength increases **linearly** from zero at the center to the outside value at the surface.



## Clicker Question

How will  $g$  change (if at all) on going from the Earth's surface to the bottom of a deep mine?  
(Assume the Earth has uniform density.)

- A.  $g$  will be a bit stronger at the bottom of the mine
- B. It will be weaker
- C. It will be the same as at the surface

## Clicker Answer

How will  $g$  change (if at all) on going from the Earth's surface to the bottom of a deep mine?

For uniform density, it will be **weaker**: the gravitational field strength varies in exactly the same way as the electric field from a solid sphere with charge uniformly distributed throughout the volume.

Note: actually the density increases with depth, so things are more complicated...

# Clicker Question

- If you could distribute charge **perfectly uniformly** throughout the volume of a solid spherical **conductor**, would it stay in place?
  - A. Yes
  - B. No

# Clicker Answer

- If you could distribute charge perfectly uniformly throughout the volume of a solid spherical conductor, would it stay that way?

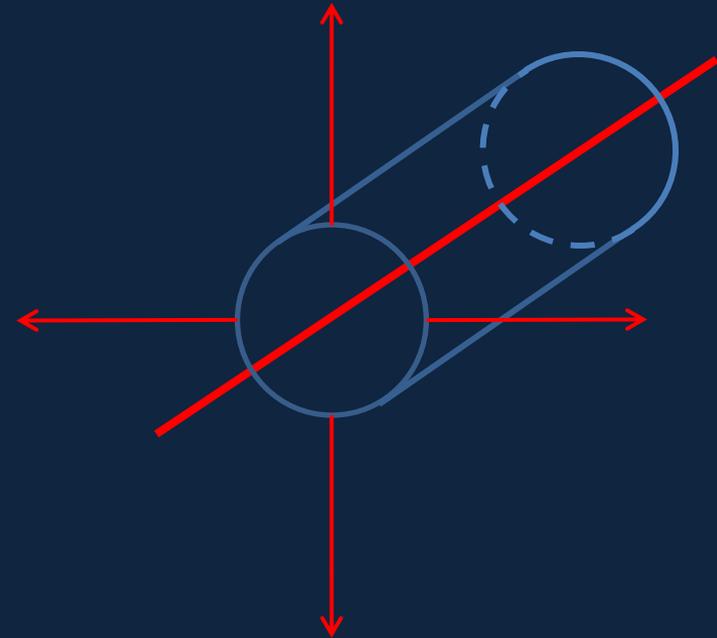
A. Yes

B. No 

Because this charge distribution gives rise to a **nonzero outward field *inside* the conductor**—the charge would therefore flow radially outwards to the surface.

# Field from a Line of Charge

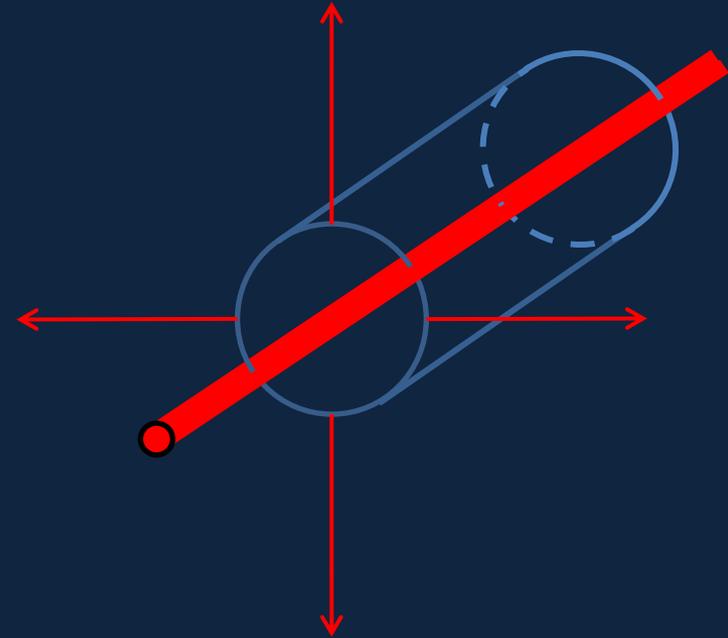
- The field is radially outward from the line, which has charge density  $\lambda$  coul/m.
- Take as gaussian surface a cylinder, radius  $r$ , axis on the line:
- The **flat ends make zero contribution** to the surface integral: the electric field vectors lie in the plane.
- For the **curved surface**:



$$\int_s \vec{E} \cdot \vec{dA} = \frac{Q}{\epsilon_0} = \frac{\lambda \ell}{\epsilon_0} = 2\pi r \ell E, \quad E = \frac{\lambda}{\epsilon_0} \frac{1}{2\pi r} = \frac{2k\lambda}{r}.$$

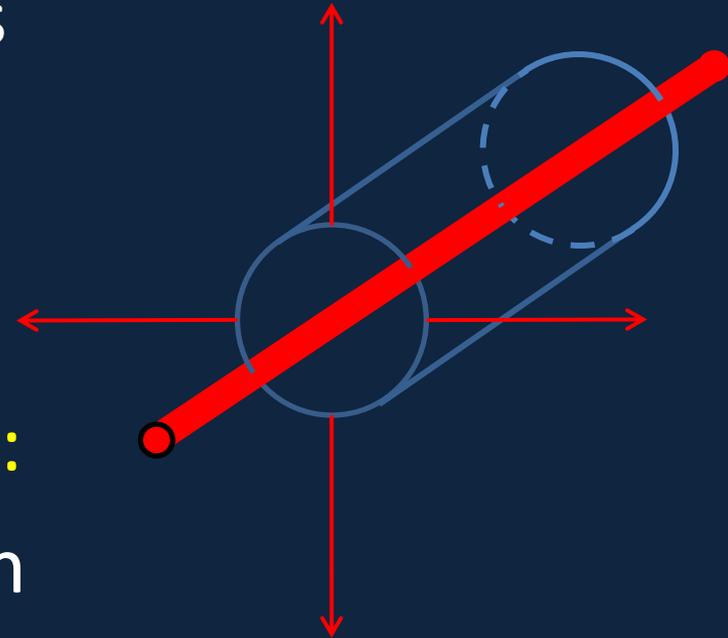
# Field from a Cylinder of Charge

- Taking a gaussian surface as shown,  $E = 2k\lambda / r$ , exactly as for a line of charge along the center.



# Clicker Question

- Suppose the central cylinder is a solid copper rod, carrying charge but with no currents anywhere.
- **The charge distribution will be:**
  - A. Uniformly distributed through the rod
  - B. Restricted to the rod's surface
  - C. Some other distribution.



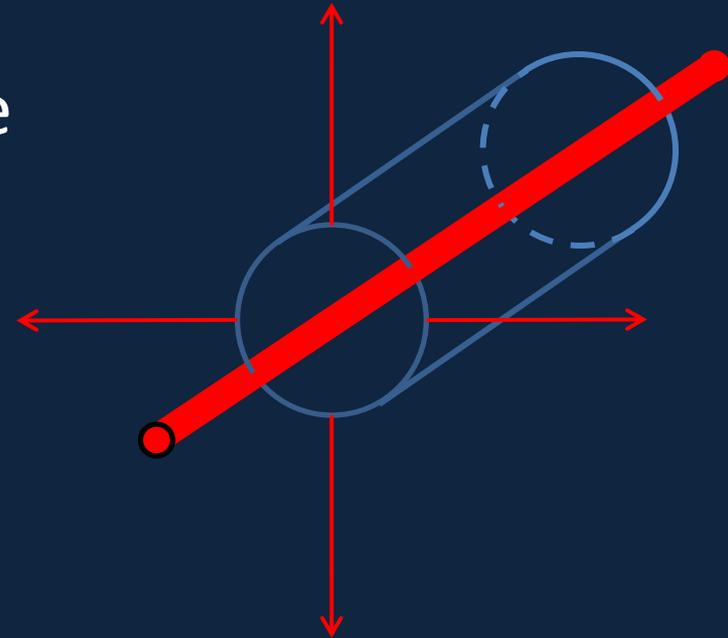
# Clicker Answer

- Suppose the central cylinder is a solid copper rod, carrying charge but with no currents anywhere.

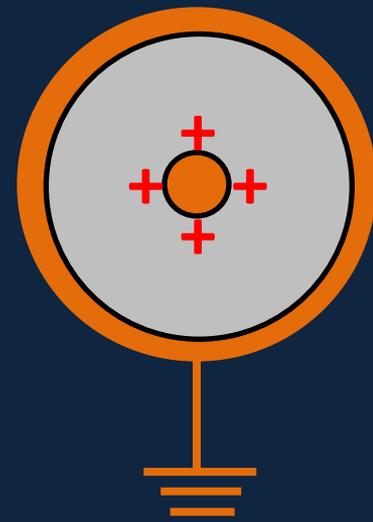
- **The charge distribution will be:**

Restricted to the rod's surface!

Just like the solid sphere, any charge inside the rod will give rise to an electric field, and therefore a current, flowing outwards.

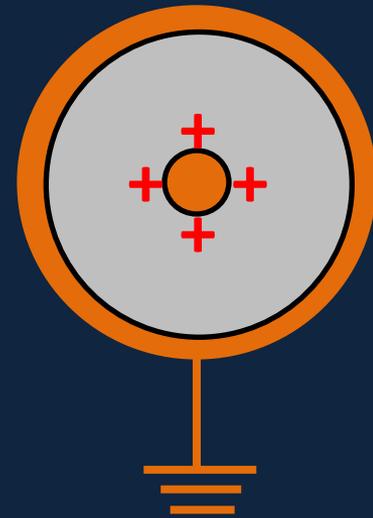


# Coaxial Cable Question



- In a coaxial cable, a central conduction cylinder is surrounded by a cylinder of insulator, and *that* is inside a hollow conducting cylinder, which is grounded here.
- If the central conductor is positively charged, the outer conducting cylinder will:
  - A. have negative charge throughout its volume
  - B. Have negative charge on its *outside* surface
  - C. Have negative charge on its *inside* surface
  - D. Have no net charge.

# Coaxial Cable Answer



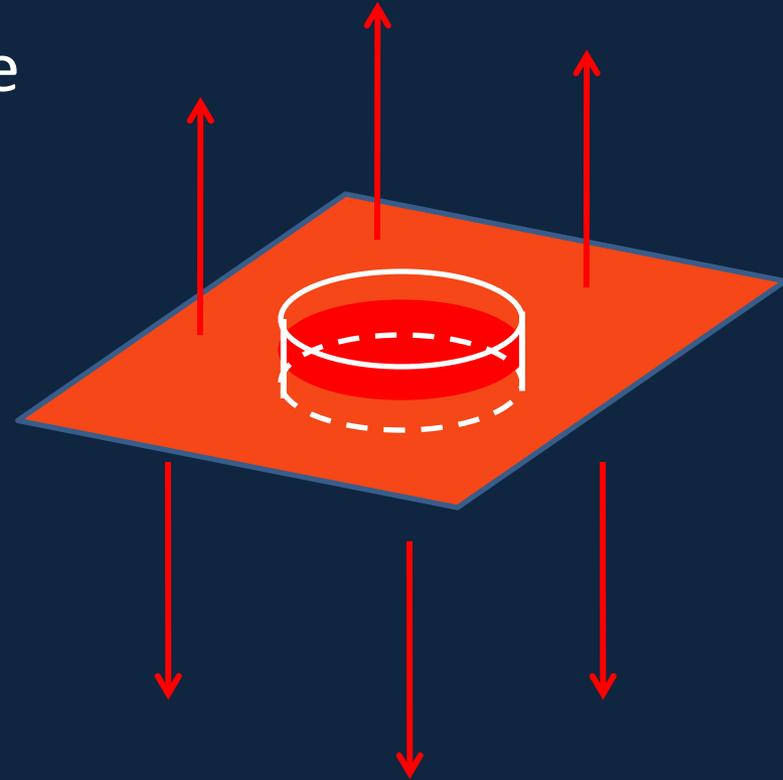
- In a coaxial cable, a central conduction cylinder is surrounded by a cylinder of insulator, and *that* is inside a hollow conducting cylinder, which is grounded here.
- **If the central conductor is positively charged, the outer conducting cylinder will:**
- Have negative charge on its *inside* surface
- The electric field lines radiating out from the inner conductor must end at the inner surface—there can be no field inside the metal of the outer cylinder.

# Uniform Sheet of Charge

- We know from symmetry that the electric field is perpendicularly outward from the plane.
- We take as gaussian surface a “pillbox”: shaped like a penny, its round faces parallel to the surface, one above and one below, area  $A$ . It contains charge (shaded red)  $Q = \sigma A$  where the charge density is  $\sigma$  C/m<sup>2</sup>.
- Gauss' theorem gives  $2AE = \frac{\sigma A}{\epsilon_0}$ , so

Both faces contribute

$$E = \frac{\sigma}{2\epsilon_0}$$



# Charge on Surface of a Conductor

- For a flat conducting surface, the electric field is perpendicularly outward, or a current would arise.
- We have a sheet of charge on the surface, so we take the same Gaussian pillbox as for the sheet of charge, but this time **there is no electric field pointing downwards into the conductor.**
- Therefore Gauss' Law gives

$$AE = \frac{\sigma A}{\epsilon_0}, \text{ so } E = \frac{\sigma}{\epsilon_0}$$

